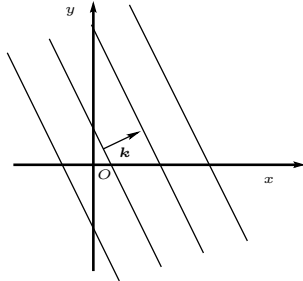


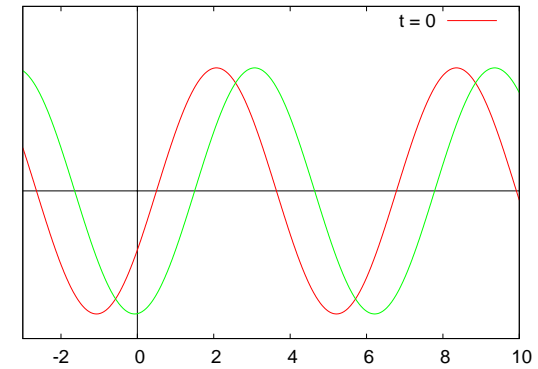
FLAT MONOCHROMATIC WAVE IN SPACE

A flat monochromatic wave in space can be given by the following vector equation:



$$E(\mathbf{r}, t) = E_0 \sin(\mathbf{k}\mathbf{r} - \omega t - \varphi) \quad (5)$$

SIN WAVE



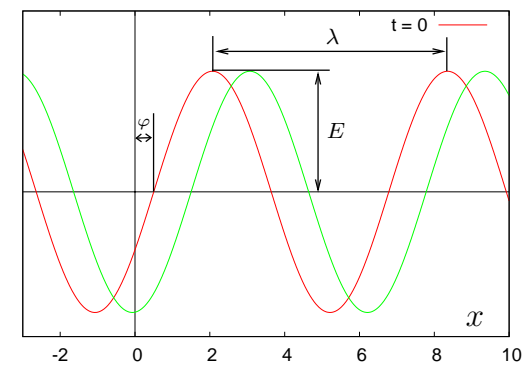
$$y = E \sin(kx - \omega t - \varphi) \quad (1)$$

SUPERPOSITION OF TWO MONOCHROMATIC WAVES

Amplitudes of two electromagnetic waves add up:

$$E = E_1 + E_2 \quad (6)$$

SIN WAVE AND ITS PARAMETERS



$$y = E \sin(kx - \omega t - \varphi) \quad (2) \quad k \stackrel{\text{def}}{=} \frac{2\pi}{\lambda} \quad (3)$$

$$\omega \stackrel{\text{def}}{=} \frac{2\pi}{T} \quad (4)$$

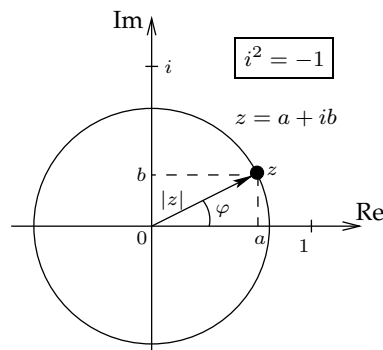
COMPLEX EXPONENTS

It is more convenient in physics to represent waves by complex exponents, using Euler's formula:

$$e^{ix} = \cos x + i \sin x \quad (10)$$

We just agree that instead of sin function, we write a complex exponent and always keep in mind that our real wave is just the imaginary part. Or real. If we care... :)

REMINESENCES: COMPLEX NUMBERS ...

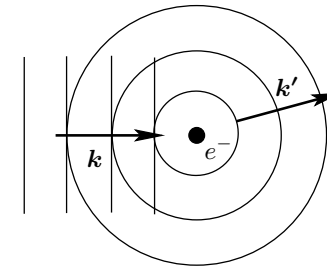


$$z = a + ib \quad |z| = \sqrt{a^2 + b^2} \quad \tan \varphi = \frac{b}{a}$$

$$z^* = a - ib \quad |z|^2 = z \cdot z^*$$

SCATTERING BY AN ELECTRON

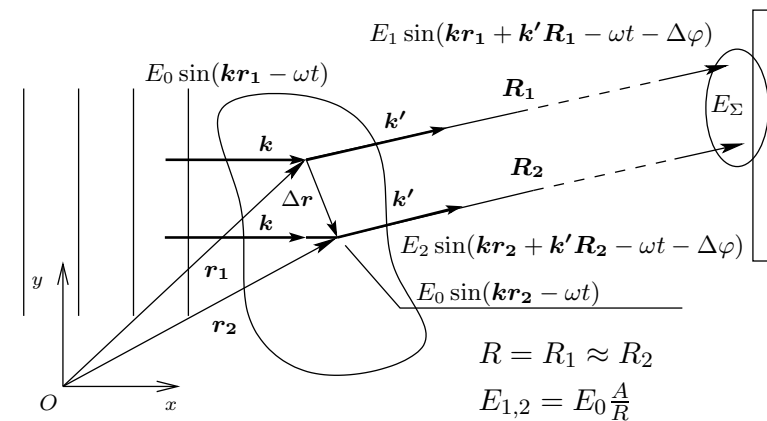
Thomson showed that:



$$E_1(\mathbf{R}, t) = E_1 \sin(\mathbf{k}' \mathbf{R} - \omega t - \Delta\varphi) \quad (7) \quad m_{p^+} \approx 1800 \cdot m_{e^-}$$

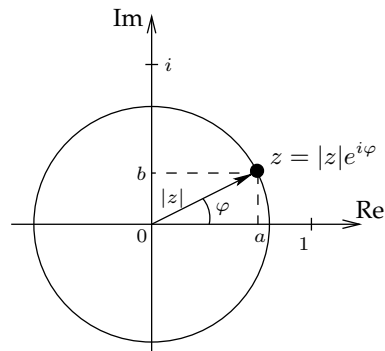
$$E_1 = E_0 \frac{1}{R} \frac{e^2}{mc^2} \sin \phi = E_0 \frac{A}{R} \quad (8) \quad I_1 \approx 2\% \cdot I_0$$

SCATTERING BY A SPECIMEN



$$E_{\Sigma} = E_1 \sin(\mathbf{k} \mathbf{r}_1 + \mathbf{k}' \mathbf{R}_1 - \omega t - \Delta\varphi) + E_2 \sin(\mathbf{k} \mathbf{r}_2 + \mathbf{k}' \mathbf{R}_2 - \omega t - \Delta\varphi) \quad (9)$$

COMPLEX NUMBER AS EXPONENTS



$$|z| = \sqrt{a^2 + b^2} \quad a = |z| \cos \varphi \quad b = |z| \sin \varphi$$

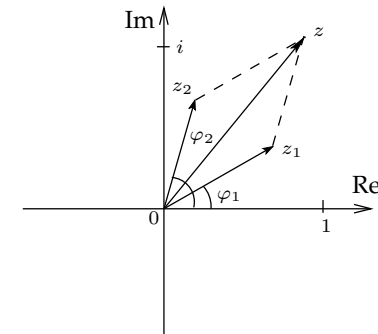
$$z = a + ib = |z| \cos \varphi + i|z| \sin \varphi = |z|(\cos \varphi + i \sin \varphi) = |z|e^{i\varphi}$$

USEFUL PROPERTIES OF EXPONENTS

$$e^a e^b = e^{a+b} \quad (13)$$

$$e^{ix} e^{i\varphi} = e^{i(x+\varphi)} \quad (14)$$

REMINESENCES: COMPLEX NUMBERS AS VECTORS ...



$$z = z_1 + z_2 = a + ib = (a_1 + a_2) + i(b_1 + b_2) \quad (11)$$

$$z_1 z_2 = (a_1 + ib_1)(a_2 + ib_2) = (a_1 a_2 - b_1 b_2) + i(a_1 b_2 + a_2 b_1) \quad (12)$$

OUTLINE OF EULER'S FORMULA PROOF...

$$e^{ix} = \cos x + i \sin x$$

$$\begin{aligned} \sin x &= x - \frac{1}{3!}x^3 + \frac{1}{5!}x^5 - \dots + (-1)^k \frac{1}{(2k+1)!}x^{2k+1} + \dots \\ i \sin x &= ix + \frac{1}{3!}(ix)^3 + \frac{1}{5!}(ix)^5 + \dots + \underbrace{(-1)^k i^{-2k-1}}_{=1} \frac{1}{(2k+1)!} (ix)^{2k+1} + \dots \\ \cos x &= 1 - \frac{1}{2!}x^2 + \frac{1}{4!}x^4 - \dots + (-1)^k \frac{1}{(2k)!}x^{2k} + \dots \\ &= 1 + \frac{1}{2!}(ix)^2 + \frac{1}{4!}(ix)^4 + \dots + \underbrace{(-1)^k i^{-2k}}_{=1} \frac{1}{(2k)!} (ix)^{2k} + \dots \\ e^x &= 1 + \frac{1}{1!}x + \frac{1}{2!}x^2 + \dots + \frac{1}{k!}x^k + \dots \\ e^{ix} &= 1 + ix + \frac{1}{2!}(ix)^2 + \frac{1}{3!}(ix)^3 + \frac{1}{4!}(ix)^4 + \dots + \frac{1}{k!}(ix)^k + \dots \end{aligned}$$

SQUARE OF THE AMPLITUDE

let

$$z = |z|e^{i\varphi}, \quad (17)$$

then

$$\begin{aligned} z \cdot z^* &= |z|e^{i\varphi} \cdot |z|e^{-i\varphi} \\ &= |z|^2 \underbrace{e^{i(\varphi-\varphi)}}_{=0} \\ &= |z|^2 \end{aligned} \quad (18)$$

REPRESENTING WAVES BY EXPONENTS

$$E \sin(\mathbf{k}\mathbf{r} - \omega t) = \text{Im}(Ee^{i(\mathbf{k}\mathbf{r} - \omega t)}) \quad (19)$$

$$E \sin(\mathbf{k}\mathbf{r} - \omega t - \varphi) = \text{Im}(Ee^{i(\mathbf{k}\mathbf{r} - \omega t - \varphi)}) \quad (20)$$

$$= \text{Im}(Ee^{-i\varphi} e^{i(\mathbf{k}\mathbf{r} - \omega t)}) \quad (21)$$

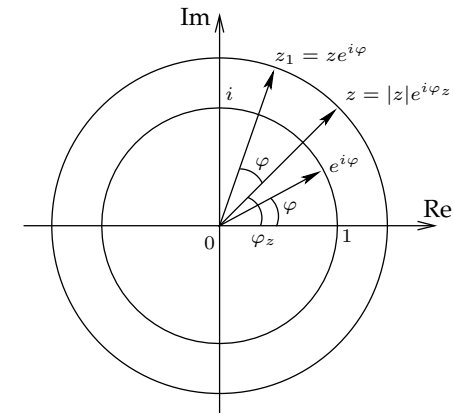
$$E_1 \sin(\mathbf{k}\mathbf{r} - \omega t - \varphi_1) +$$

$$E_2 \sin(\mathbf{k}\mathbf{r} - \omega t - \varphi_2) = \text{Im}(E_1 e^{i(\mathbf{k}\mathbf{r} - \omega t - \varphi_1)} + E_2 e^{i(\mathbf{k}\mathbf{r} - \omega t - \varphi_2)})$$

$$= \text{Im}((E_1 e^{-i\varphi_1} + E_2 e^{-i\varphi_2}) e^{i(\mathbf{k}\mathbf{r} - \omega t)}) \quad (22)$$

$$(23)$$

PHASE SHIFT



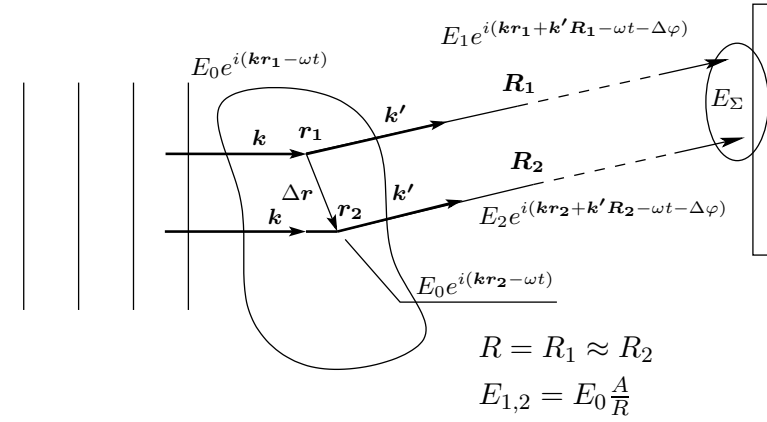
$$z_1 = ze^{i\varphi_z} = |z|e^{i\varphi_z} e^{i\varphi} = |z|e^{i(\varphi_z + \varphi)} \quad (15)$$

ADDING TWO EXPONENTS

$$\begin{aligned} e^{ix} + e^{iy} &= e^{ix}(1 + e^{i(y-x)}) \\ &= e^{ix} e^{i\frac{y-x}{2}} (e^{-i\frac{y-x}{2}} + e^{i\frac{y-x}{2}}) \\ &= e^{i2\frac{x}{2}} e^{i\frac{y-x}{2}} \left(\cos \frac{y-x}{2} - i \sin \frac{y-x}{2} + \cos \frac{y-x}{2} + i \sin \frac{y-x}{2} \right) \\ &= e^{i\frac{x+y}{2}} \cdot 2 \cos \frac{y-x}{2} \end{aligned} \quad (16)$$

SCATTERING BY A SPECIMEN REVISITED

exponential form:

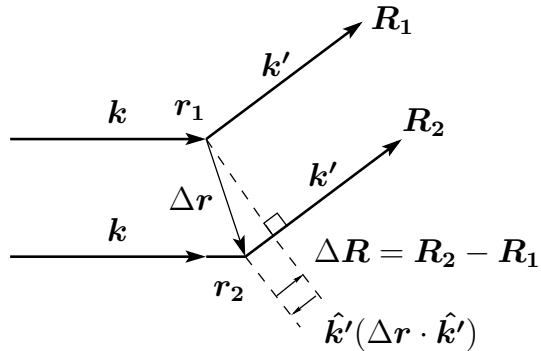


$$E_\Sigma = E_1 e^{i(kr_1 + k'R_1 - \omega t - \Delta\varphi)} + E_2 e^{i(kr_2 + k'R_2 - \omega t - \Delta\varphi)} \quad (24)$$

ADDING TWO MONOCHROMATIC WAVES (2)

$$\begin{aligned}
 E_\Sigma &= \underbrace{E'_0 e^{i(kr_1 + k'R_1)}}_{=E'_0} [1 + e^{i(kr_2 - kr_1 + k'R_2 - k'R_1)}] \\
 &= E'_0 [1 + e^{i(k\Delta r + k'\Delta R)}]
 \end{aligned} \quad (28)$$

PROJECTION OF Δr



$$\begin{aligned}
 \hat{k}' &= \frac{\mathbf{k}'}{k'} & \Delta R &= -\hat{k}'(\Delta r \cdot \hat{k}') \\
 \mathbf{k}' \Delta R &= -\mathbf{k}' \cdot \frac{\mathbf{k}'}{k'} (\Delta r \cdot \frac{\mathbf{k}'}{k'}) = -\frac{(\mathbf{k}' \cdot \mathbf{k}')}{k'^2} (\Delta r \cdot \mathbf{k}') = -\mathbf{k}' \Delta r
 \end{aligned}$$

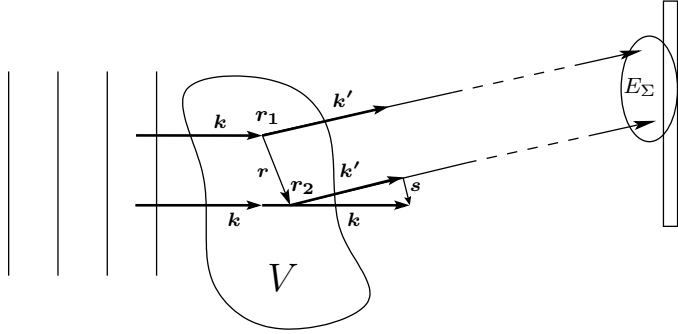
ADDING TWO MONOCHROMATIC WAVES

$$\begin{aligned}
 E_1(\mathbf{R}_1, t) &= E_1 e^{i(kr_1 + k'R_1 - \omega t - \Delta\varphi)} \\
 &= \underbrace{E_0 \frac{A}{R} e^{-i\Delta\varphi} e^{-i\omega t} e^{i(kr_1 + k'R_1)}}_{\stackrel{\text{def}}{=} E'_0} \\
 &= E'_0 e^{i(kr_1 + k'R_1)}
 \end{aligned} \quad (25)$$

$$E_2(\mathbf{R}_2, t) = E_0 e^{i(kr_2 + k'R_2)} \quad (26)$$

$$\begin{aligned}
 E_\Sigma &= E'_0 e^{i(kr_1 + k'R_1)} + E_0 e^{i(kr_2 + k'R_2)} \\
 &= E'_0 [e^{i(kr_1 + k'R_1)} + e^{i(kr_2 + k'R_2)}] \\
 &= E'_0 e^{i(kr_1 + k'R_1)} [1 + e^{i(kr_2 - kr_1 + k'R_2 - k'R_1)}]
 \end{aligned} \quad (27)$$

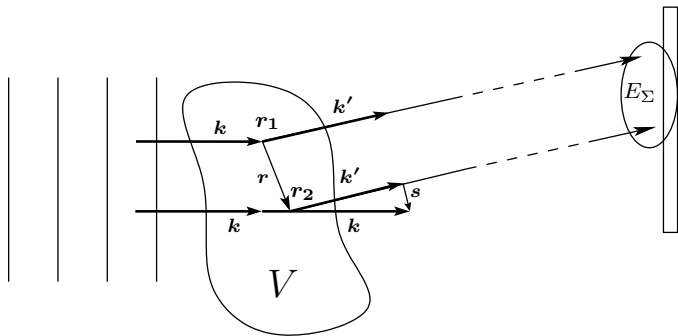
FOURIER TRANSFORM



$$E_\Sigma = \lim_{\substack{N \rightarrow \infty \\ \Delta V_k \rightarrow 0}} E_0'' \sum_{k=1}^N \rho_k \Delta V_k e^{i\mathbf{r}_k \cdot \mathbf{s}} \quad (31)$$

$$= E_0'' \int_V \rho(\mathbf{r}) e^{i\mathbf{r} \cdot \mathbf{s}} dV \quad (32)$$

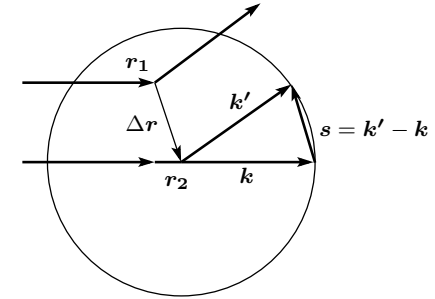
FOURIER TRANSFORM (2)



$$\mathbf{S} = \frac{\mathbf{s}}{2\pi} \quad (33)$$

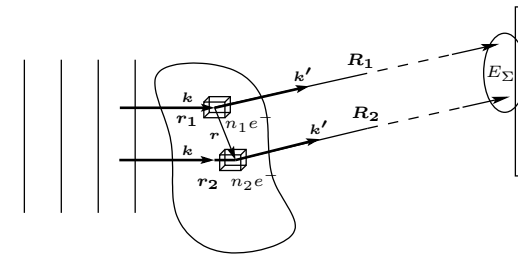
$$E_\Sigma = E_0'' \int_V \rho(\mathbf{r}) e^{2\pi i \mathbf{r} \cdot \mathbf{S}} dV \quad (34)$$

ADDING TWO MONOCHROMATIC WAVES (3)



$$\begin{aligned} E_\Sigma &= E_0'' [1 + e^{i(\mathbf{k}\Delta\mathbf{r} - \mathbf{k}'\Delta\mathbf{r})}] \\ &= E_0'' [1 + e^{i(\Delta\mathbf{r}(\mathbf{k} - \mathbf{k}'))}] \\ &= E_0'' [1 + e^{i\Delta\mathbf{r}\mathbf{s}}] \\ &= E_0'' [1 \cdot e^{i\Delta\mathbf{r}_0\mathbf{s}} + e^{i\Delta\mathbf{r}\mathbf{s}}] \end{aligned} \quad (29)$$

ADDING TWO MONOCHROMATIC WAVES (4)



$$\begin{aligned} E_{1+2} &= E_0'' [n_1 e^{i\mathbf{r}_0 \cdot \mathbf{s}} + n_2 e^{i\mathbf{r} \cdot \mathbf{s}}] \\ E_\Sigma &= E_0'' \sum_{k=1}^N n_k e^{i\mathbf{r}_k \cdot \mathbf{s}} \\ &= E_0'' \sum_{k=1}^N \rho_k \Delta V_k e^{i\mathbf{r}_k \cdot \mathbf{s}} \end{aligned} \quad (30)$$

PROPERTIES OF THE FOURIER TRANSFORM

- Linearity

$$\mathcal{F}[\alpha\rho_1 + \beta\rho_2] = \alpha\mathcal{F}[\rho_1] + \beta\mathcal{F}[\rho_2] \quad (39)$$

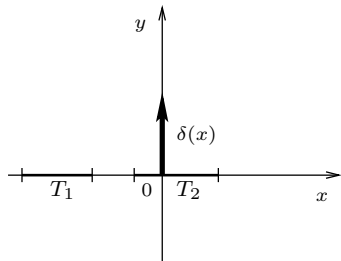
- Shift

$$\mathcal{F}[\rho(\mathbf{r} - \Delta\mathbf{r})] = e^{2\pi i(\Delta\mathbf{r}\mathbf{S})} \mathcal{F}[\rho(\mathbf{r})] \quad (40)$$

- Rotation

$$\begin{aligned} \mathbf{r}' &= ||R||\mathbf{r} \\ \rho'(\mathbf{r}) &\stackrel{def}{=} \rho(\mathbf{r}') = \rho(||R||\mathbf{r}) \\ \mathcal{F}[\rho(||R||\mathbf{r})] &= F(||R||\mathbf{S}) \end{aligned} \quad (41)$$

DIRAC DELTA FUNCTION



$$\int_{T_1} \delta(x)f(x)dx \stackrel{def}{=} 0 \quad (42)$$

$$\int_{T_2} \delta(x)f(x)dx \stackrel{def}{=} f(0) \quad (43)$$

$$\int_{-\infty}^{+\infty} \delta(x)f(x)dx \stackrel{def}{=} f(0) \quad (44)$$

STRUCTURE FACTOR

$$E_{\Sigma} = E_0'' \underbrace{\int_V \rho(\mathbf{r})e^{2\pi i\mathbf{r}\mathbf{S}} dV}_{F(\mathbf{S})} \quad (35)$$

$$\boxed{\mathcal{F}[\rho(\cdot)](\mathbf{S}) = F(\mathbf{S}) = \int_V \rho(\mathbf{r})e^{2\pi i\mathbf{r}\mathbf{S}} dV} \quad (36)$$

INVERSE FOURIER TRANSFORM

let

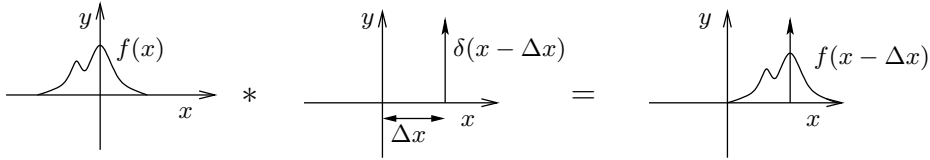
$$F(\mathbf{S}) = \int_{-\infty}^{+\infty} \rho(\mathbf{r})e^{2\pi i\mathbf{r}\mathbf{S}} dV \quad (37)$$

then

$$\rho(\mathbf{r}) = \int_{-\infty}^{+\infty} F(\mathbf{S})e^{-2\pi i\mathbf{r}\mathbf{S}} d\mathbf{S} \quad (38)$$

PROPERTIES OF THE DELTA FUNCTION

CONVOLUTION WITH DELTA FUNCTION



$$f(x) * \delta(x - \Delta x) = \int_{-\infty}^{+\infty} f(v) \delta(x - v - \Delta x) dv = f(x - \Delta x) \quad (52)$$

$$\delta(-x) = \delta(x) \quad (45)$$

$$\int_{-\infty}^{+\infty} \delta(x) dx = 1 \quad (46)$$

$$\int_{-\infty}^{+\infty} \delta(x - x_0) f(x) dx = f(x_0) \quad (47)$$

$$\mathcal{F}[\delta] = \int_{-\infty}^{+\infty} \delta(x) e^{2\pi i x S} dx = e^0 = 1 \quad (48)$$

$$\mathcal{F}^{-1}[1] = \int_{-\infty}^{+\infty} e^{-2\pi i x S} dS = \delta(x) \quad (49)$$

$$\mathcal{F}[1] = \int_{-\infty}^{+\infty} e^{2\pi i x S} dx = \delta(S) \quad (50)$$

FOURIER TRANSFORM OF THE CONVOLUTION

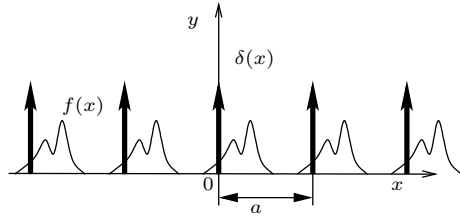
$$\begin{aligned} \mathcal{F}[f * g] &= \int_{-\infty}^{+\infty} \left(\int_{-\infty}^{+\infty} f(v) g(x - v) dv \right) e^{2\pi i S x} dx \\ &= \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} f(v) g(x - v) e^{2\pi i S(x-v)} e^{2\pi i S v} dv dx \\ &= \int_{-\infty}^{+\infty} f(v) e^{2\pi i S v} dv \cdot \int_{-\infty}^{+\infty} g(x - v) e^{2\pi i S(x-v)} d(x - v) \end{aligned} \quad (53)$$

$$= \mathcal{F}[f] \cdot \mathcal{F}[g] \quad (54)$$

CONVOLUTION

$$(f * g)(x) \stackrel{\text{def}}{=} \int_{-\infty}^{+\infty} f(v) g(x - v) dv \quad (51)$$

CRYSTAL FOURIER TRANSFORM

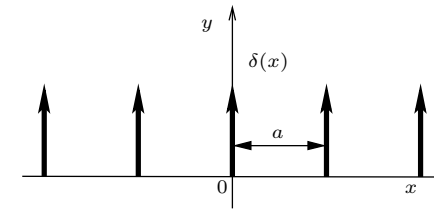


$$\mathcal{F}[f * L] = \mathcal{F}[f] \cdot \mathcal{F}[L] \quad (57)$$

$$= \mathcal{F}\left[\sum_{n=-\infty}^{+\infty} f(x - an)\right] \quad (58)$$

$$= \sum_{n=-\infty}^{+\infty} \mathcal{F}[f(x - an)] \quad (59)$$

LATTICE FUNCTION



$$L(x) = \sum_{n=-\infty}^{+\infty} \delta(x - an) \quad (55)$$

CRYSTAL FOURIER TRANSFORM (2)

$$\mathcal{F}[f * L] = \mathcal{F}[f] \cdot \mathcal{F}[L] \quad (60)$$

$$= F(S) \cdot \mathcal{F}[L] \quad (61)$$

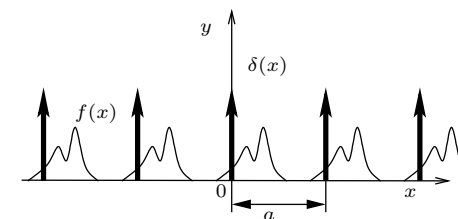
$$= \mathcal{F}\left[\sum_{n=-\infty}^{+\infty} f(x - an)\right] \quad (62)$$

$$= \sum_{n=-\infty}^{+\infty} \mathcal{F}[f(x - an)] \quad (63)$$

$$= \sum_{n=-\infty}^{+\infty} F(S) e^{2\pi i S a n} \quad (64)$$

$$= F(S) \sum_{n=-\infty}^{+\infty} e^{2\pi i S a n} \quad (65)$$

CRYSTAL FUNCTION



$$f * L = \sum_{n=-\infty}^{+\infty} f(x - an) \quad (56)$$

CRYSTAL FOURIER TRANSFORM (4)

$$c_n = \frac{1}{a} F(n/a) \quad (73)$$

$$\sum_{n=-\infty}^{\infty} f(x - an) = \sum_{n=-\infty}^{\infty} c_n e^{-2\pi i \cdot n \frac{x}{a}} \quad (74)$$

$$= \frac{1}{a} \sum_{n=-\infty}^{\infty} F(n/a) e^{-2\pi i \cdot n \frac{x}{a}} \quad (75)$$

$$\sum_{n=-\infty}^{\infty} \delta(x - an) = \frac{1}{a} \sum_{n=-\infty}^{\infty} e^{-2\pi i \cdot n \frac{x}{a}} \quad (76)$$

LATTICE FOURIER TRANSFORM

$$\sum_{n=-\infty}^{\infty} \delta(x - an) = \frac{1}{a} \sum_{n=-\infty}^{\infty} e^{2\pi i \cdot n \frac{x}{a}} \quad (77)$$

$$\sum_{n=-\infty}^{\infty} \delta(x - n/a) = a \sum_{n=-\infty}^{\infty} e^{-2\pi i \cdot n a x} \quad (78)$$

CRYSTAL FOURIER TRANSFORM (3)

$$\sum_{n=-\infty}^{\infty} f(x - an) \stackrel{\text{def}}{=} S(x) \quad (66)$$

$$= \sum_{n=-\infty}^{\infty} c_n e^{-2\pi i \cdot n \frac{x}{a}} \quad (67)$$

$$c_n = \frac{1}{a} \int_{-a/2}^{a/2} S(x) e^{2\pi i \cdot n \frac{x}{a}} dx \quad (68)$$

$$= \frac{1}{a} \int_{-a/2}^{a/2} \sum_{n=-\infty}^{\infty} f(x - an) e^{2\pi i \cdot n \frac{x}{a}} dx \quad (69)$$

$$= \frac{1}{a} \sum_{n=-\infty}^{\infty} \int_{-a/2}^{a/2} f(x - an) e^{2\pi i \cdot n \frac{x}{a}} dx \quad (70)$$

$$= \frac{1}{a} \int_{-\infty}^{\infty} f(x) e^{2\pi i \cdot x \frac{n}{a}} dx \quad (71)$$

$$= \frac{1}{a} F(n/a) \quad (72)$$

RECIPROCAL LATTICE, 3D CASE

$$L(x) = \sum_{m,n,p=-\infty}^{+\infty} \delta(\mathbf{r} - \mathbf{a}m - \mathbf{b}n - \mathbf{c}p) \quad (87)$$

$$L^*(\mathbf{S}) = \mathcal{F}[L] = \frac{1}{V} \sum_{h,k,l=-\infty}^{+\infty} \delta(\mathbf{S} - h\mathbf{a}^* - k\mathbf{b}^* - l\mathbf{c}^*) \quad (88)$$

$$\mathbf{a}^* = \frac{[\mathbf{b} \times \mathbf{c}]}{V}, \quad \mathbf{b}^* = \frac{[\mathbf{c} \times \mathbf{a}]}{V}, \quad \mathbf{c}^* = \frac{[\mathbf{a} \times \mathbf{b}]}{V} \quad (89)$$

$$V = (\mathbf{a} \cdot \mathbf{b} \times \mathbf{c}) = (\mathbf{a} \cdot [\mathbf{b} \times \mathbf{c}]) \quad (90)$$

LAUE CONDITIONS

$$\mathcal{F}[\rho_{cryst}] = \mathcal{F}[\rho * L] = \mathcal{F}[\rho] \cdot \mathcal{F}[L] \quad (91)$$

$$= F(\mathbf{S}) \cdot L^*(\mathbf{S}) \quad (92)$$

$$= F(\mathbf{S}) \cdot \frac{1}{V} \sum_{h,k,l=-\infty}^{+\infty} \delta(\mathbf{S} - h\mathbf{a}^* - k\mathbf{b}^* - l\mathbf{c}^*) \quad (93)$$

We see that the transformant of the crystal is non-vanishing only for a certain values of \mathbf{S}

$$(\mathbf{S} \cdot \mathbf{a}) = h, \quad (\mathbf{S} \cdot \mathbf{b}) = k, \quad (\mathbf{S} \cdot \mathbf{c}) = l \quad (94)$$

$$h, k, l \in \mathbb{Z} \quad (95)$$

$$\mathcal{F}[L] = \mathcal{F}\left[\sum_{n=-\infty}^{\infty} \delta(x - an)\right] \quad (79)$$

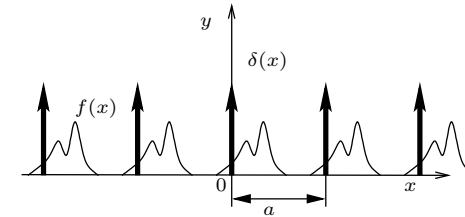
$$= \int_{-\infty}^{\infty} \sum_{n=-\infty}^{\infty} \delta(x - an) e^{2\pi i x S} dx \quad (80)$$

$$= \sum_{n=-\infty}^{\infty} \int_{-\infty}^{\infty} \delta(x - an) e^{2\pi i x S} dx \quad (81)$$

$$= \sum_{n=-\infty}^{\infty} e^{2\pi i a n S} \quad (82)$$

$$= \frac{1}{a} \sum_{n=-\infty}^{\infty} \delta(S - n/a) \quad (83)$$

RECIPROCAL LATTICE, 1D CASE



$$L(x) = \sum_{n=-\infty}^{+\infty} \delta(x - an) \quad (84)$$

$$\mathcal{F}[f * L] = \mathcal{F}[f] \cdot \mathcal{F}[L] \quad (85)$$

$$= F(S) \cdot \frac{1}{a} \sum_{n=-\infty}^{\infty} \delta(S - n/a) \quad (86)$$

EVARD CONSTRUCT AND EVARD SPHERE

